

The epitaxy of AlN(11 $\bar{2}2$)-layers on GaN(11-22)/m-Al₂O₃

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The epitaxial layers of AlN(11 $\bar{2}2$) with a thickness of 3.9 μm were grown on a GaN(11 $\bar{2}2$)/m-Al₂O₃ template by hydride vapour-phase epitaxy (HVPE). The template consisted of GaN(11 $\bar{2}2$)- and buffer AlN s with thicknesses of 2.7 μm and 0.6 μm , respectively, grown on a sapphire substrate of orientation (10 $\bar{1}0$). It is shown that the full width at half maximum (FWHM) for the diffraction peaks of GaN(11 $\bar{2}2$)/m-Al₂O₃ and AlN(11 $\bar{2}2$) are 30 and 20 arc minutes, respectively. It was found that the epitaxy of AlN(11 $\bar{2}2$) on the template leads to an increase in the size of crystal blocks in the AlN(11 $\bar{2}2$) layer. It is assumed that the improvement in the quality of the AlN(11 $\bar{2}2$) layer occurs due to its predominant growth in the tangential direction due to the relatively low lattice difference at the AlN(11 $\bar{2}2$)/GaN(11 $\bar{2}2$) heterogeneous boundary.

Keywords: semipolar AlN(11 $\bar{2}2$), hydride vapour-phase epitaxy.

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Aluminum nitride (AlN) has great potential for application in optoelectronic devices operating in the deep ultraviolet range and in high-power and high-frequency electronic devices, since it has a wide band gap and fitting properties [1]. However, strong spontaneous and piezoelectric polarization in the [0001] direction of aluminum nitride reduces the overlap of the electron and hole wave functions and exerts a negative influence on the electron recombination efficiency. This obstacle may be cleared through the use of non-polar or semi-polar layers. However, it is extremely difficult to grow high-quality semi-polar AlN due to the low migration activity of aluminum atoms on the surface, which inevitably leads to a columnar growth mode and a rough surface. Several methods for improving the quality of crystalline polar and semi-polar epitaxial AlN layers, such as high-temperature growth [2], sapphire nanostructuring [3], or preliminary nitridation of the sapphire substrate [4], have been developed. It is rather difficult to grow layers with a smooth surface due to the strong anisotropy of non-polar and semi-polar AlN [5].

The use of AlN(10 $\bar{1}1$), AlN(10 $\bar{1}3$), and AlN(11 $\bar{2}2$) layers in device structures for suppression of the influence of internal electric fields is a promising trend in optoelectronics [6]. Numerous studies focused on improving the quality of semi-polar AlN on sapphire have been published over the last two decades, and the possibility of synthesis of layers with different semi-polar surface orientations has been demonstrated [7]. The growth of high-quality AlN layers by metal-organic chemical vapor deposition (MOCVD) and hydride vapor-phase epitaxy (HVPE) with magnetron sputtering of AlN on sapphire as a buffer layer has recently been reported in [8].

In the present study, we propose a new approach to HVPE synthesis of AlN(11 $\bar{2}2$) layers on a sapphire substrate

with the (10 $\bar{1}0$) surface orientation with micrometer-sized GaN(11 $\bar{2}2$) layers used as a template.

Let us note first that AlN(11 $\bar{2}2$) layers may be formed on an m-Al₂O₃ substrate by heating in a gas environment with excess ammonia. This is made possible by the formation of nanofacets on m-sapphire at high temperature [9]. AlN(11 $\bar{2}2$) layers were synthesized as follows. Heteroepitaxy was performed on sapphire substrates of the (10 $\bar{1}0$) orientation cleaned in advance in the standard way. A buffer layer of AlN with a thickness of 0.6 μm was grown first in an argon atmosphere, which was followed by the synthesis of a 2.7- μm -thick GaN layer. At the last stage, the main layer of AlN with a thickness of 3.9 μm was formed (Fig. 1). The flow rates of HCl and NH₃ were 1.7 and 2.4 l/min, respectively. The temperature of epitaxy of AlN

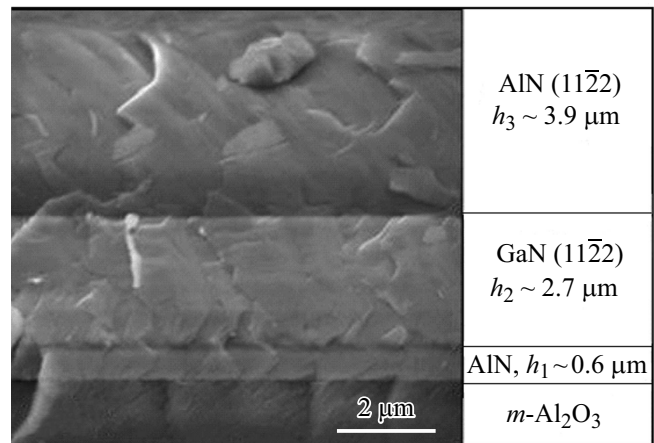


Figure 1. SEM image of the cleaved surface of an AlN(11 $\bar{2}2$) layer with template GaN(11 $\bar{2}2$)/m-Al₂O₃.

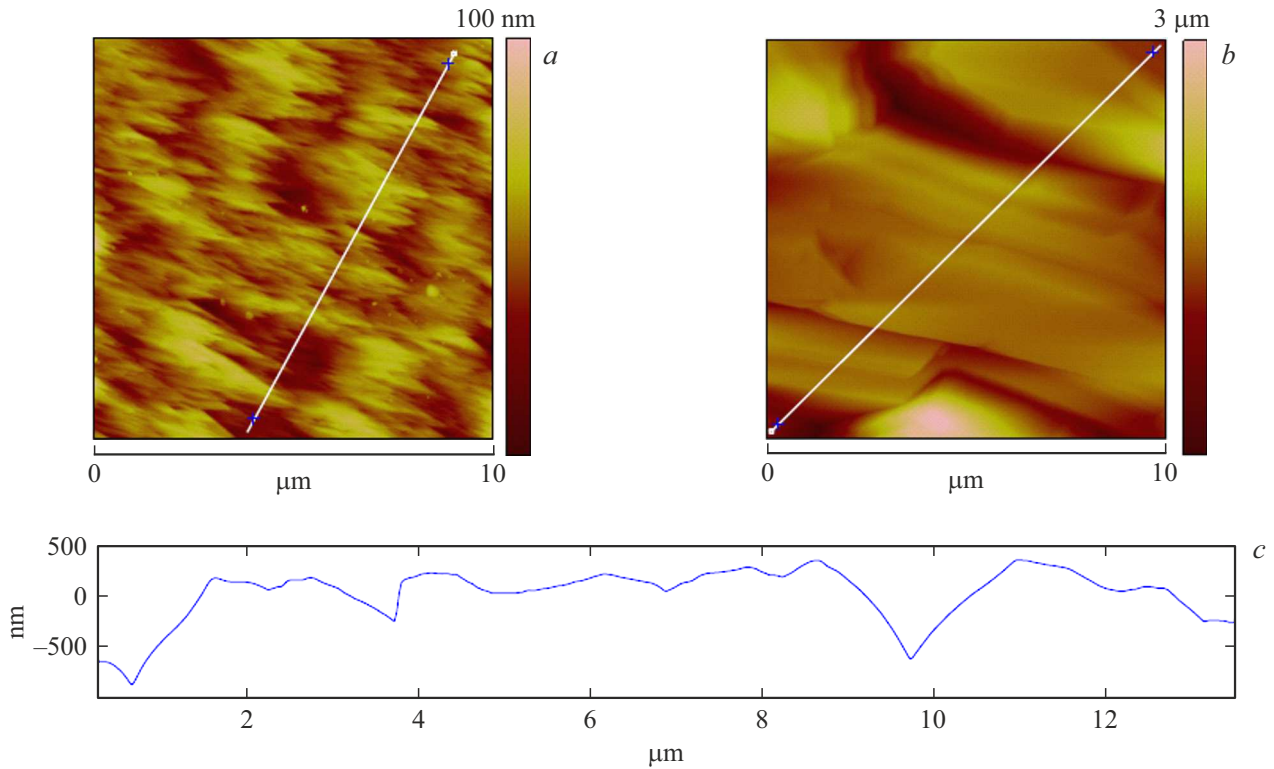


Figure 2. AFM images (*a*, *b*) and profile (*c*) of the surface of GaN(11 $\bar{2}2$)/ m -Al₂O₃ (*a*) and AlN(11 $\bar{2}2$)/GaN(11 $\bar{2}2$)/ m -Al₂O₃ (*b*, *c*) structures.

Table 1. FWHM of X-ray diffraction ω_θ , root-mean-square value (RMS), and average roughness R_a for a semi-polar AlN(11 $\bar{2}2$) layer and a GaN(11 $\bar{2}2$)/ m -Al₂O₃ template with scanning over 1×1 and $10 \times 10 \mu\text{m}$ areas

Substrate	ω_θ , arcmin	RMS, nm		R_a , nm	
		1×1	10×10	1×1	10×10
GaN(11 $\bar{2}2$)/ m -Al ₂ O ₃	30	1.65	11.2	1.26	8.87
AlN(11 $\bar{2}2$)/GaN(11 $\bar{2}2$)/ m -Al ₂ O ₃	20	2.35	116	1.38	85.5

and GaN layers was close to 1000 and 1050 °C, respectively. The approximate time of epitaxial layer growth was 15 min.

The structural and optical characteristics of GaN and AlN layers of these structures were determined using X-ray diffractometry and optical, scanning electron (SEM), and atomic force (AFM) microscopy. In X-ray measurements, rocking curves were recorded in a double-crystal diffraction arrangement in reflections (0002) and (11 $\bar{2}2$) of CuK α_1 radiation using a triple-crystal X-ray spectrometer.

Nitridation of the m -Al₂O₃ substrate at $T = 1000$ °C provided an opportunity to grow a low-quality AlN buffer layer consisting of blocks, some of which have the AlN(11 $\bar{2}2$) orientation. Partial orientation of the buffer layer in the semi-polar direction allowed us to grow GaN(11 $\bar{2}2$) and AlN(11 $\bar{2}2$) layers.

X-ray diffraction measurements revealed that the full width at half maximum (FWHM) of X-ray rocking curves for the GaN(11 $\bar{2}2$) and AlN(11 $\bar{2}2$) diffraction peaks is 30 and 20 arcmin, respectively.

AFM imaging of GaN(11 $\bar{2}2$) layers and, after further growth, AlN(11 $\bar{2}2$) layers demonstrated that the surface of both layers consists of elongated domains (Fig. 2). It was found that small domains $2 \times 4 \mu\text{m}$ in size on the surface of the GaN(11 $\bar{2}2$) layer merged in the course of further growth of AlN, dwindling in number and increasing in size to $8 \times 20 \mu\text{m}$, and higher-quality AlN(11 $\bar{2}2$) layers were formed. Compared to GaN(11 $\bar{2}2$), the root-mean-square (RMS) and average (R_a) roughness values of the AlN(11 $\bar{2}2$) layer surface increased slightly when scanning over an area of $1 \times 1 \mu\text{m}$ and differed significantly at $10 \times 10 \mu\text{m}$ (Table 1).

As is known, the AlN orientation on m -plane sapphire in gas-phase epitaxy may change either to (10 $\bar{1}3$) or to (11 $\bar{2}2$) depending on the concentration of NH₃ in the gas phase and the temperature [9]. This is attributable to the possible formation of r -plane nanofacets on m -plane sapphire in the case of its nitridation and at a relatively high temperature, which lead to growth of the (11 $\bar{2}2$) plane.

Table 2. Lattice parameter a at room temperature [10] and relative mismatch $\Delta a/a$ at the heteroboundary

Lattice parameter a , Å					Lattice parameter mismatch $\Delta a/a$	
$m\text{-Al}_2\text{O}_3$	GaN	AlN	GaN(11 $\bar{2}2$)	AlN(11 $\bar{2}2$)	GaN(11 $\bar{2}2$)/ $m\text{-Al}_2\text{O}_3$	AlN(11 $\bar{2}2$)/GaN(11 $\bar{2}2$)
4.785	3.189	3.112	5.517	5.390	0.128	-0.024

According to the results of X-ray diffraction studies, the semi-polar orientation of the buffer layer cannot be established, apparently due to the polycrystalline nature of the buffer layer resulting from a significant difference in lattice parameters of $m\text{-Al}_2\text{O}_3$ and AlN (Table 2).

We associate the effect of growth of AlN(11 $\bar{2}2$) blocks with the difference in magnitude of the lattice parameter mismatch at the GaN(11 $\bar{2}2$)/ $m\text{-Al}_2\text{O}_3$ and AlN(11 $\bar{2}2$)/GaN(11 $\bar{2}2$) heteroboundaries (Table 2). The relative magnitude of mismatch at the heteroboundary may be estimated as $\Delta a/a = (a_1 - a_2)/a_2$, where a_1 and a_2 are the lattice parameters of two adjacent layers at the heteroboundary (Table 2) [10]. We believe that the large lattice mismatch at the AlN(11 $\bar{2}2$)/ $m\text{-Al}_2\text{O}_3$ and GaN(11 $\bar{2}2$)/ $m\text{-Al}_2\text{O}_3$ heteroboundaries leads, first, to the formation of a fine-grained polycrystalline AlN buffer layer and, second, to the formation of a layer of semi-polar gallium nitride with block sizes of about $2\mu\text{m}$ in the $[1\bar{1}00]$ direction (see Fig. 2, *a*). When the AlN(11 $\bar{2}2$) layer is formed on the GaN(11 $\bar{2}2$)/ $m\text{-Al}_2\text{O}_3$ template, the lattice mismatch in this direction is significantly smaller (Table 2). This leads to the growth of a layer with larger blocks by HVPE in an argon atmosphere (Figs. 2, *b*, *c*). The FWHM for diffraction peaks of the template and the AlN(11 $\bar{2}2$) layer was found to be 30 and 20 arcmin, respectively, which is close to similar diffraction data for the AlN(11 $\bar{2}2$) layers grown on (10 $\bar{1}0$) sapphire substrates by MOCVD [11].

The results demonstrate the potential of the method for fabrication of stripe templates for micrometer-sized semi-polar optoelectronic structures. It was hypothesized that the AlN(11 $\bar{2}2$) layer quality improves due to a relatively slight lattice mismatch at the AlN(11 $\bar{2}2$)/GaN(11 $\bar{2}2$) heteroboundary.

Conflict of interest

The authors declare that they have no conflict of interest.

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