#### 10,11

# Lattice dynamics and phase transitions in a crystal Ba<sub>2</sub>ZnTeO<sub>6</sub>

#### © V.I. Zinenko, A.S. Shinkorenko¶

Kirensky Institute of Physics of Federal Research Center KSC of the Siberian Branch of RAS, Krasnoyarsk, Russia <sup>¶</sup> E-mail: shas@iph.krasn.ru

Received July 20, 2022 Revised July 20, 2022 Accepted July 21, 2022

The lattice dynamics of Ba<sub>2</sub>ZnTeO<sub>6</sub> double perovskite in three phases has been studied. Vibrational spectra and dynamic charges are calculated. In the rhomobohedral phase  $R\bar{3}m$ , soft modes were found; after their condensation, the C2/m phase was obtained, which agrees with the experiment. An assessment was made of the dependence of the enthalpy of different phases on pressure; it showed that the monoclinic phase is more favorable at 0 pressure, but as the pressure increases to 2.9 GPa, the cubic phase  $Fm\bar{3}m$  becomes advantageous.

Keywords: Lattice dynamics, phase transitions, double perovskites.

DOI: 10.21883/PSS.2022.12.54395.445

## 1. Introduction

Double perovskites with the general formula  $A_2BB'O_6$ are being studied for several decades due to their structural phase transitions and diversity of physical properties, which depend both on chemical composition of A, B and B'cations and on different types and degree of cation ordering in the structure [1]. Most of these compounds in the highly symmetrical cubic phase (with the space group of symmetry  $Fm\bar{3}m$ ) have the elpasolite structure where octahedra  $BO_6$ and  $B'O_6$  are linked by oxygen at the octahedron vertexes. Most double perovskites with the elpasolite structure (with few exceptions), regardless of the degree of cation ordering in the cubic phase, have instabilities related to lattice vibration modes either in the center or at the grating boundary point X of the Brillouin zone in a face-centered cubic lattice, the eigen vectors of these "soft" modes correspond to "turns" of BO6 octahedron. In rarer cases, with a large ionic radius of the A cations, double perovskites can have other structure types, in particular, with hexagonal or a rhombohedral symmetry [2-4]. The BO<sub>6</sub> and B'O<sub>6</sub> octahedra in such structures have common vertexes, edges and faces. Such double perovskites with common edges and faces of BO<sub>6</sub> and B'O<sub>6</sub> octahedra are less well understood both from the viewpoint of their physical properties and structure instability.

Double perovskite Ba<sub>2</sub>ZnTeO<sub>6</sub> has rhombohedral symmetry with space group  $R\bar{3}m$  and with two molecules in the unit cell.

The results of the study of Raman and IR spectra of Ba<sub>2</sub>ZnTeO<sub>6</sub> were published in [5,6]. The authors of these papers have found a structural phase transition  $(R\bar{3}m \rightarrow C2/m)$  at 140 K, which is related to the "soft" vibration mode. The authors have also found the crystal lattice vibration frequencies and classified them. The authors of [7] studied the behavior of Ba<sub>2</sub>ZnTeO<sub>6</sub> under hydrostatic pressure and found that at higher pressures  $Ba_2ZnTeO_6$  has the elpasolite structure with symmetry  $Fm\bar{3}m$ .

The aim of the present paper is to calculate the lattice vibration frequencies and to study the phase transitions in the  $Ba_2ZnTeO_6$  compound using the ab initio calculation methods implemented in the Crystal software package [8,9].

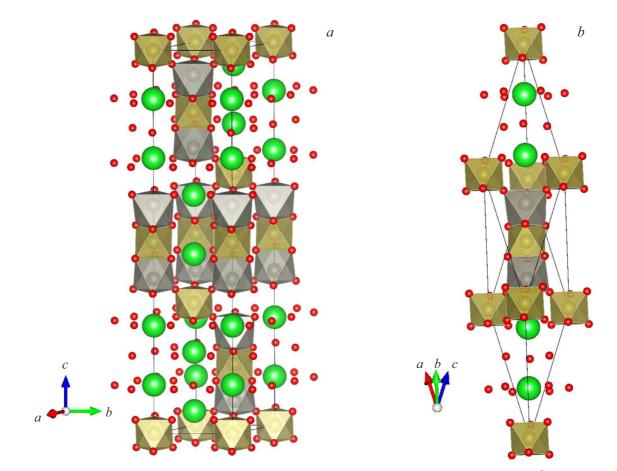
# 2. Structural characteristics and calculation method

The crystal structure of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound is shown in Fig. 1. Several types of octahedra can be distinguished in the structure. The first type is TeO<sub>6</sub> octahedra located in layers in plane *ab* and linked to ZnO<sub>6</sub> octahedra by oxygen at the octahedron vertexes. The second type is TeO<sub>6</sub> octahedra located in columns along axis *c* and linked by faces with ZnO<sub>6</sub> octahedra. The third type is ZnO<sub>6</sub> octahedra located in columns along axis *c*; on one side they have a common face with a TeO<sub>6</sub> octahedron of the second type, on the other side their vertexes are linked to TeO<sub>6</sub> octahedra of the first type. A shift of axis *c* of the Zn atom from the octahedron center should also be noted.

The calculations in the present paper were performed using the CRYSTAL package. The hybrid B3LYP method was used in the DFT calculations [10]. The chosen bases are (HAYWSC-31(2df)G) for the barium atom, (Zn\_pob\_TZVP\_rev2) for the zinc atom, (Te\_POB\_TZVP\_2018) for the tellurium atom and (O\_8-411d11G\_valenzano\_2006) for the oxygen atom [11].

### 3. Results

We have calculated the equilibrium lattice parameters and the coordinates of atoms in the structure with the symmetry group  $R\bar{3}m$ , which are given in Table 1 together



**Figure 1.** Crystal structure of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the rhombohedral phase with the space group  $R\bar{3}m$ : a — unit cell and b — primitive cell.

with the experimental data. An expansion of a vibrational representation at the Brillouin zone center in indecomposable representations of point group  $D_{3d}$  has the form  $\Gamma = 7A_{1g} + 2A_{2g} + 9E_g + 2A_{1u} + 9A_{2u} + 11E_u$ . The calculated vibration frequencies of atoms at the Brillouin zone center are given in Table 2 together with the experimental data [5,6]). The values of dynamic ion charges are given in Table 3.

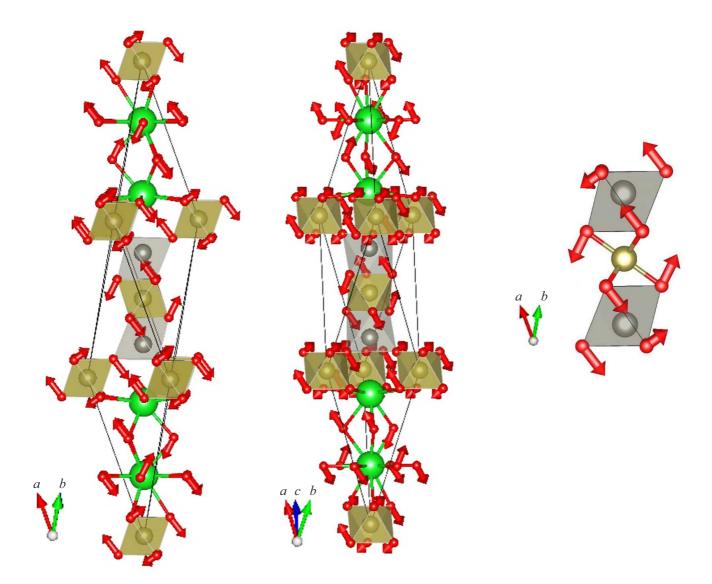
It can be seen from Table 2 that the calculated vibration frequencies mainly agree well with the experimental values, but there are some discrepancies both between the experimental and calculated values of vibration frequencies and between the experimental data. In particular, both experimental spectra have a high-frequency vibration  $E_g$  with the frequency of 767 cm<sup>-1</sup>, but vibrations with this frequency are not presented in the calculated spectrum.

Two soft vibration modes are present in the calculated vibration frequencies of the BZTO compound in the structure with symmetry group  $R\bar{3}m$ . One doubly degenerate mode  $E_g$  with the frequency of  $-67 \text{ cm}^{-1}$  and non-degenerate mode  $A_{2g}$  with the frequency of  $-53 \text{ cm}^{-1}$ . The observed structural phase transition  $R\bar{3}m \rightarrow C2/m$  is related to the unstable mode  $E_g$ .

**Table 1.** Calculated and experimental (\*) [7] lattice parameters and coordinates of atoms of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the rhombohedral phase with space group  $R\bar{3}m$ 

Lattice parameters, Å	5.9 (5.8	c 29.13 (28.69*)	
Atom	x	У	z
Ba <sub>1</sub>	0	0	0.6274
$Ba_2$	0	0	0.7815
Zn	0	0	-0.0963
Te <sub>1</sub>	0	0	0.5
Te <sub>2</sub>	0	0	0
$O_1$	0.8199	0.1800	0.2931
$O_2$	0.1783	0.8216	0.1272

Let us consider atomic displacements in the eigen vector of the imaginary vibration mode  $E_g$ , which are shown by arrows in Fig. 2, *a*. It can be seen from this figure that displacements of oxygen ions at the vertexes of the TeO<sub>6</sub>



**Figure 2.** Eigen vector of one component of the dual mode  $E_g$  for Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the rhombohedral phase with space group  $R\bar{3}m$ .

octahedra both of the first and the second type correspond to a "turn" of this octahedra, as in double perovskites with the elpasolite structure where the octahedra are linked by the vertexes. However,  $ZnO_6$  octahedra, as distinct from the elpasolite structure, are distorted in thiscase.

A distortion of the structure of the highly symmetrical phase  $R\bar{3}m$  along the eigen vectors of duel mode  $E_g$ causes the formation of a structure with symmetry group C2/m. It should be noted that a distortion of the structure along one component or along two equal components of the eigen vector of mode  $E_g$  causes the formation of the same structure (An expansion of the thermodynamic potential in the order parameter, which is converted by the presentation  $E_g$ , contains a single quartic invariant). The total crystal energy vs. the amplitude of atomic displacement along one component of the eigen vector is shown in Fig. 3. The optimized lattice parameters and coordinates of atoms of the structure with group C2/m are given in Table 4. The experimental data is also given there. It should be noted that there is a disagreement with the experimental data on the lattice parameter, namely parameter a — the difference between the calculated and the experimental value is  $\approx 10\%$ .

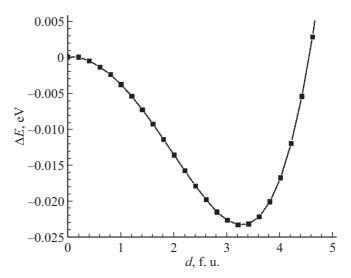
The calculated vibration frequencies at the center of the Brillouin zone of the distorted structure are given in Table 5. An expansion in indecomposable representations at the center of the Brillouin zone for symmetry group C2/mis  $\Gamma = 16A_g + 11B_g + 13A_u + 20B_u$ . The experimental frequencies are also given there. Good agreement with the experiment can be seen. A soft vibration mode  $B_g$  with the frequency of  $-57.3 \text{ cm}^{-1}$  is present in the calculated vibration frequencies of the BZTO compound in the structure with symmetry group C2/m. The symmetry

	$E_g$			$A_{1g}$			$A_{2g}$	
Calc.	Exp.*	Exp.**	Calc.	Exp.*	Exp.**	Calc.	Exp.*	Exp.**
-67.2	28.8	31	84.5	86.0	87.0	-53.7		
79.3	121.4	103	116.2	104.3	110.0	130.9	185	185.0
95.6	382.7	120	219.1	109.7	405.0			
122.7	395.5	153	392.3	142.5	470.0			
205.0	406.2	382	440.9	153.7	689.0			
365.3	572.0	394	665.3	471.3	736.0			
384.2	616.0	573	760.0	756.1				
555.5	690.8	616						
629.2	767.7	766						
E	'g	$A_1$	lg	P	$\Lambda_{2g}$			
Calc.	Exp.*	Calc.	Exp.*	Calc.	Exp.*			
0.0	67.1	200.8	231.0	0.0	99			
52.9	109.6	233.3	331.5	101.3	141.6			
116.9	124.5			106.3	182.2			
122.3	135.2			173.1	192.1			
208.2	259.5			220.2	200.6			
234.8	312.1			357.0	437.5			
282.4	361.2			438.4	475.4			
345.4	388.4			650.7	755.3			
367.0	619.2			756.9				
595.1	660.4							
664.6	1	1	1	1	1	1		

**Table 2.** Calculated and experimental  $[5^*, 6^{**}]$  vibration frequencies  $(cm^{-1})$  of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the rhombohedral phase with space group  $R\bar{3}m$ 

**Table 3.**Calculated dynamic charges of the  $Ba_2ZnTeO_6$ compound in the rhombohedral phase with space group  $R\bar{3}m$ 

Atom	$Z_{xx}$	$Z_{yy}$	$Z_{zz}$
Ba <sub>1</sub>	2.8	2.8	2.7
Ba <sub>2</sub>	2.8	2.8	2.7
Ba <sub>3</sub>	2.9	2.9	2.71
Ba <sub>4</sub>	2.9	2.9	2.71
$Zn_1$	2.37	2.37	2.71
$Zn_2$	2.37	2.37	2.71
Te <sub>1</sub>	5.05	5.05	4.92
Te <sub>1</sub>	4.29	4.29	5.23
$O_1$	-2.63	-1.79	-1.98
$O_2$	-2.63	-1.79	-1.98
O3	-2.63	-1.98	-1.79
$O_4$	-2.63	-1.98	-1.79
$O_5$	-2.63	-1.79	-1.98
$O_6$	-2.63	-1.79	-1.98
$O_7$	-3.14	-1.73	-1.67
$O_8$	-3.14	-1.73	-1.67
O9	-3.14	-1.67	-1.73
$O_{10}$	-3.14	-1.67	-1.73
O <sub>11</sub>	-3.14	-1.73	-1.67
O <sub>12</sub>	-3.14	-1.73	-1.67



**Figure 3.** Total energy of BZTO crystal in the monoclinic phase vs. shift amplitude.

decreases to group  $P\bar{1}$  in case of a crystal structure distortion along the eigen vector of the given vibration. As per the literature data, such a transition is not observed in the BZTO crystal, but there is experimental data [12]

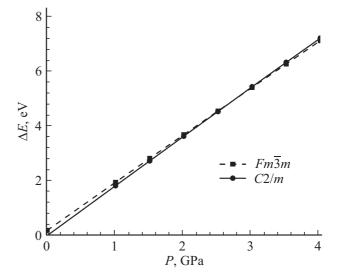
	а	b	с
Lattice parameters, Å	10.185 9.0228*	5.9186 5.8236*	10.289 10.086*
	90	109.08 109.36*	90
Atom	x	у	z
Ba <sub>1</sub>	0.1258	0	0.3853
$Ba_2$	0.2870	0	0.8453
Zn	0.3975	0	0.2102
Te <sub>1</sub>	0	0	0
Te <sub>2</sub>	0	0.5	0.5
$O_1$	0.3815	0.2303	0.3821
$O_2$	0.6081	0	0.3741
$O_3$	0.5587	0.2658	0.8956
$O_4$	0.8154	0	0.8610

**Table 4.** Calculated and experimental [10] lattice parameters and coordinates of atoms of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the monoclinic phase with space group C2/m

that the isomorphic Ba<sub>2</sub>CuTeO<sub>6</sub> compound undergoes a  $C2/m \rightarrow P\bar{1}$  transition at 290 K.

The values of dynamic ion charges in phase C2/m are given in Table 6.

As already stated in the introduction, the authors of [7] found a phase transition to the elpasolite structure with space symmetry  $Fm\bar{3}m$  in the BZTO crystal under hydrostatic pressure. We optimized the BZTO structure in terms of lattice parameters and atom coordinates in phase  $Fm\bar{3}m$ . The results are given in Table 7. The vibration frequencies at



**Figure 4.** Enthalpy  $\Delta H$  of the BZTO compound in C2/m and  $Fm\bar{3}m$  phases vs. applied pressure.

the center of the Brillouin zone and dynamic Born charges were calculated for this optimized structure. The results are given in Table 7, 8. It should be emphasized that the BZTO crystal in the elpasolite structure has no unstable vibration modes, i.e. this structure remains stable at all temperatures.

Enthalpy  $\Delta H = E_{\text{full}} - PV$  of the BZTO compound in phases C2/m and  $Fm\bar{3}m$  vs. applied pressure is shown in Fig. 4. The zero is the energy at 0 pressure of the structure with symmetry group C2/m, all energies were standardized to the formula unit. Enthalpy vs. pressure is not given for the phase with structure  $R\bar{3}m$ , because this

**Table 5.** Calculated and experimental [5] vibration frequencies  $\omega$ , cm<sup>-1</sup> of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the monoclinic phase with space group C2/m

Calc.		IR		Calc.	Calc.		Raman		Calc.
$A_u$	$A_u$	$A_u + B_u$	$B_u$	B <sub>u</sub>	$A_g$	$A_g$	$A_g + B_g$	Bg	$B_g$
Au           38.4           119.6           135.1           184.4           209.3           224.9           235.1           291.5           344.5           369.5           593.2           647.5	331.5	$\begin{array}{c} A_u + B_u \\ 63 \\ 71.7 \\ 105 \\ 109.6 \\ 122 \\ 128.2 \\ 133.7 \\ 139.2 \\ 251 \\ 266 \\ 310 \\ 314.5 \\ 358.9 \\ 386.9 \\ 613 \\ 621.3 \end{array}$	100.4 143 183.7 193.4 201 438.5 476	<i>B<sub>u</sub></i> 68.3 90.6 99.1 122.1 135.8 179.7 208.0 220.4 279.6 300.8 343.8 362.6 367.3 449.1 589.9 640.8	$\begin{array}{c} A_g \\ 72.8 \\ 77.2 \\ 96.9 \\ 103.5 \\ 124.6 \\ 153.0 \\ 206.3 \\ 232.1 \\ 367.8 \\ 384.2 \\ 414.6 \\ 450.8 \\ 555.5 \\ 609.3 \\ 660.7 \\ 737.8 \end{array}$	$\begin{array}{c} A_g \\ 92.6 \\ 99.4 \\ 108.7 \\ 140.6 \\ 166 \\ 469.8 \\ 756.1 \end{array}$	$\begin{array}{c} A_g + B_g \\ 37.3 \\ 50.7 \\ 119.1 \\ 125.9 \\ 379.2 \\ 382.2 \\ 397.9 \\ 399.9 \\ 404.6 \\ 406.6 \\ 574.2 \\ 577.6 \\ 615.9 \\ 621.8 \\ 693.6 \\ 695.7 \end{array}$	184.8	$B_g$ -57.3 47.8 89.2 113.9 126.9 144.6 195.9 358.4 380.9 551.6 609.2
		660 664.3	761.5	647.8 734.2		773.3	770.9		

Atom	$Z_{xx}$	$Z_{yy}$	$Z_{zz}$
Ba <sub>1</sub>	2.89	2.61	2.76
Ba <sub>2</sub>	2.89	2.61	2.76
Ba <sub>3</sub>	2.73	2.92	2.95
Ba <sub>4</sub>	2.73	2.92	2.95
$Zn_1$	2.74	2.37	2.37
$Zn_2$	2.74	2.37	2.37
$Te_1$	4.94	4.94	5.18
Te <sub>2</sub>	4.31	5.19	4.27
$O_1$	-2.57	-1.79	-2.01
$O_2$	-2.57	-1.79	-2.01
$O_3$	-2.68	-1.95	-1.79
$O_4$	-2.68	-1.95	-1.79
$O_5$	-2.57	-1.79	-2.01
$O_6$	-2.57	-1.79	-2.01
$O_7$	-3.16	-1.73	-1.64
$O_8$	-3.16	-1.73	-1.64
O9	-3.01	-1.77	-1.74
$O_{10}$	-3.01	-1.77	-1.74
O <sub>11</sub>	-3.16	-1.73	-1.64
O <sub>12</sub>	-3.16	-1.73	-1.64

**Table 6.**Calculated dynamic charges of the  $Ba_2ZnTeO_6$ compound in the monoclinic phase with space group C2/m

**Table 7.** Calculated and experimental [7] lattice parameters and coordinates of atoms of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the cubic phase with space group  $Fm\bar{3}m$ 

Lattice	a 8.2518 (8.1479*)			
parameter, Å				
Atom	x	у	z	
Ba Zn Te O	0.25 0 0.5 0.2636 0.2520*	0.25 0 0.5 0	0.25 0 0.5 0	

**Table 8.** Calculated vibration frequencies of the Ba<sub>2</sub>ZnTeO<sub>6</sub> crystal in the cubic phase with structure  $Fm\bar{3}m$ 

Sim.	$\omega$ , cm <sup>-1</sup>
$F_{1u}$	0
$F_{1g}$	55.6
$F_{2g}$	120.5
$F_{1u}$	121.2
$F_{1u}$	158.5
$F_{2u}$	236
$F_{1u}$	332.8
	372.6
$F_{2g} \ E_g$	446.6
$F_{1u}$	555.4
$A_g$	648.9

in the cubic phase	e with structure	e Fm3m	
Atom	$Z_{xx}$	$Z_{yy}$	$Z_{zz}$
Ba <sub>1</sub>	2.73	2.73	2.73
Ba <sub>2</sub>	2.73	2.73	2.7.3

Table 9. Calculated dynamic charges of the Ba<sub>2</sub>ZnTeO<sub>6</sub> crystal

$Ba_2$	2.73	2.73	2.73
Zn	2.96	2.96	2.96
Te	4.76	4.76	4.76
$O_1$	-1.69	-3.19	-1.69
$O_2$	-1.69	-1.69	-3.19
$O_3$	-1.69	-1.69	-3.19
$O_4$	-3.19	-1.69	-1.69
$O_5$	-1.69	-3.19	-1.69
$O_6$	-3.19	-1.69	-1.69
	•	-	

phase does not become more favorable than the others at all the studied pressure values (calculations are performed at zero temperature). As seen in Fig. 4, the phase with structure  $Fm\bar{3}m$  becomes more favorable than the phase with structure C2/m at 2.9 GPa. The obtained data agrees with the experiment [7], where this phase transition was found at the pressure of 4.0 GPa.

# 4. Conclusion

In the present paper we have calculated the vibration frequencies at the center of the Brillouin zone of the Ba<sub>2</sub>ZnTeO<sub>6</sub> compound in the rhombohedral, monoclinic and cubic phases. Imaginary vibration modes were obtained in the spectrum of the calculated frequencies for the rhombohedral phase: doubly degenerate mode  $E_g$  with the frequency of  $-67 \text{ cm}^{-1}$  and non-degenerate mode  $A_{2g}$  with the frequency of  $-53 \text{ cm}^{-1}$ . The structural phase transition  $R\bar{3}m \rightarrow C2/m$  is related to the unstable mode  $E_g$ . An experimentally observed monoclinic phase of BZTO was obtained after distortion along one component of mode  $E_g$ . A soft vibration mode Bg with the frequency of  $-57.3 \text{ cm}^{-1}$  was found in the spectrum of the calculated frequencies of the monoclinic phase. After its condensation, a structure with symmetry group  $P\bar{1}$  was obtained.

The enthalpy of the different phases vs. pressure was estimated. It was found that the cubic phase with structure  $Fm\bar{3}m$  becomes more favorable than the monoclinic phase with structure C2/m at 2.9 GPa. The obtained data agrees with the experiment [7], where this phase transition was found at the pressure of 4.0 GPa.

#### Funding

The calculations were performed using the computer resources "Integrated simulation and processing of data of mega-class scientific research units "of National Research Centre "Kurchatov Institute" (http://ckp.urcki.ru).

### **Conflict of interest**

The authors declare that they have no conflict of interest.

# References

- [1] S. Vasala, M. Karppinen. Prog. Solid State Chem. 43, 1 (2015).
- [2] P. Von Köhl, U. Müller, and D. Reinen, Z. Anorg. Allg. Chem. 392, 124 (1972).
- [3] G.M. Keith, C.A. Kirk, K. Sarma, N.McN. Alford, E.J. Cussen, M.J. Rosseinsky, D.C. Sinclair. Chem. Mater. 16, 2007 (2004).
- [4] S.-F. Wang, Y.-F. Hsu, H.-S. Huang, Y.-J. Liu. Ceram. Int. 37, 1327 (2011).
- [5] R.L. Moreira, R.P.S.M. Lobo, S.L.L.M. Ramos, M.T. Sebastian, F.M. Matinaga, A. Righi, A. Dias. Phys. Rev. Mater. 2, 054406 (2018).
- [6] S. Badola, B. Ghosh, G. Sunil, S. Saha. arXiv:2110.12430
- [7] T. Aoba, T. Tiittanen, H. Suematsu, M. Karppinen. J. Solid State Chem. 233, 492 (2016).
- [8] R. Dovesi, R. Orlando, A. Erba, C.M. Zicovich-Wilson,
  B. Civalleri, S. Casassa, L. Maschio, M. Ferrabone,
  M. De La Pierre, P. D'Arco, Y. Noel, M. Causa, M. Rerat,
  B. Kirtman. Int. J. Quantum Chem. 114, 1287 (2014).
- [9] R. Dovesi, V.R. Saunders, C. Roetti, R. Orlando, C.M. Zicovich-Wilson, F. Pascale, B. Civalleri, K. Doll, N.M. Harrison, I.J. Bush, P. D'Arco, M. Llunell, M. Causá, Y. Noël. CRYSTAL14 User's Manual. University of Torino, Torino (2014).
- [10] D. Becke. Density-functional thermochemistry. III The role of exact exchange. J. Chem. Phys. 98, 5648 (1993).
- [11] https://www.crystal.unito.it/basis-sets.php
- [12] A.S. Gibbs, A. Yamamoto, A.N. Yaresko, K.S. Knight, H. Yasuoka, M. Majumder, M. Baenitz, P.J. Saines, J.R. Hester, D. Hashizume, A. Kondo, K. Kindo, H. Takagi. Phys. Rev. B 95, 104428 (2017).